

# Study of Higher Order Dispersion Effects on Soliton Propagation In Dispersion Shifted Fiber

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**ABSTRACT :** In this paper I studied higher order dispersion ( $\beta_3$ ) effects on soliton propagation in dispersion shifted fiber (DSF) with solving numerically the Nonlinear Schrödinger Equation (NLSE). Initial pulse have hyperbolic-secant shape and power that give a fundamental soliton, fiber loss is compensated by pre-emphasis method. I used Symmetrized Split Step Fourier (SSSF) method to solve the NLSE and compare soliton propagation in DSF from result of NLSE with  $\beta_3$  is included and  $\beta_3$  is negligible by measure root mean square (rms) pulse width of soliton that determines the bit rate of soliton communication systems.

**KEYWORDS :** Higher order dispersion ( $\beta_3$ ), soliton propagation, pre-emphasis method, dispersion shifted fiber (DSF), Nonlinear Schrödinger Equation (NLSE), Symmetrized Split Step Fourier (SSSF) method.

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## 1 INTRODUCTION

The word *soliton* was coined by Krusal and Zabusky in 1965 to describe the particle like properties of pulse envelope in dispersive nonlinear media; under certain condition the envelope not only propagates undistorted but survives collisions just as particles do [1].

In 1973 Hasegawa and Tappert suggested the possibility of soliton propagation in optical fiber in their paper on *Applied Physic. Letter. 23, 142 (1973)* [2]. Soliton in optical fibers were first observed by L.F. Mollenauer *et al.* in an experiments that used a mode-locked color-center laser to obtain short ( $T_{FWHM} \approx 7$  ps,  $FWHM = Full\ Width\ at\ Half\ Maximum$ ) optical pulse near 1.55  $\mu\text{m}$ , a wavelength close to which the fiber loss is minimum [2].

DSF is fiber that zero dispersion wavelength shifted to near 1,55  $\mu\text{m}$ , here fiber loss is minimum. In DSF when operating wavelength is 1.55  $\mu\text{m}$ , GVD (Group Velocity Delay) parameter ( $\beta_2$ ) is close to zero so higher order dispersion ( $\beta_3$ ) must be considered [1].

The propagation of optical pulse envelope for pulse width  $\geq 0.1$ ps in a nonlinear fiber included higher order dispersion and maintain its polarization is goverend by Nonlinear Shrödinger Equation (NLSE), can be written as [1]

$$\frac{\partial A}{\partial z} + \frac{a}{2}A + \frac{i}{2}b_2 \frac{\partial^2 A}{\partial T^2} - \frac{1}{6}b_3 \frac{\partial^3 A}{\partial T^3} = ig |A|^2 A \quad (1)$$

where :

$A$  = Pulse envelope (slowly varying variable).

$g$  = Nonlinearity parameter.

$T = t - z/v_g$ , where  $T$  = time in a reference frame moving with pulse;  $t$  = physical time,  $z$  = the direction of pulse propagation along fiber axis and  $v_g$  = the group velocity at the center wavelength. Eq. (1) can be written in normalized form as:

$$\frac{\partial U}{\partial x} + \left( L_D \cdot \frac{a}{2} \right) U + \text{sgn}(b_2) \frac{i}{2} \frac{\partial^2 U}{\partial t^2} - \text{sgn}(b_3) \frac{1}{6} b_3 \frac{L_D}{L_D} \frac{\partial^3 U}{\partial t^3} = iN^2 |U|^2 U \quad (2)$$

where :  $t = \frac{T}{T_0}$  ;  $L_D = \frac{T_0^2}{|b_2|}$  ;  $L_D' = \frac{T_0^3}{|b_3|}$  ;  $L_N = \frac{1}{gP_{avg}}$  ;  $x = \frac{z}{L_D}$  ;  $N^2 = \frac{L_D}{L_N}$  ;  $U = \frac{A}{\sqrt{P_0}}$  and

$\text{sgn}(\beta_2)$  (or  $\text{sgn}(\beta_3)$ ) takes values +1 and -1 depending on whether  $\beta_2$  (or  $\beta_3$ ) is positive or negative [1].  $U$ ,  $a$ ,  $P_0$ ,  $\xi$ ,  $T_0$ ,  $L_D$ ,  $L_D'$  and  $L_N$  is normalized pulse envelope, fiber loss, peak input pulse power, input pulse width, normalized distance, dispersion length, higher order dispersion length and nonlinear length. Soliton exist only for integer values of  $N$ , soliton corresponding to  $N = 1$  is called fundamental soliton [1].

## 2 RMS PULSE WIDTH AND PRE-EMPHASIS METHOD

### 2.1 RMS Pulse Width

In a dispersive nonlinear fiber, the pulse shape at the fiber output can deviate considerably from the input pulse shape and can have much more complicated forms. RMS pulse width is a proper measure of the pulse width that propagates in nonlinear fiber. *RMS pulse width* defined by [4] :

$$s = \left[ \langle t^2 \rangle - \langle t \rangle^2 \right]^{\frac{1}{2}} \quad (3)$$

where :

$$\langle t^n \rangle = \frac{\int_{-\infty}^{\infty} t^n |U(z,t)|^2 dt}{\int_{-\infty}^{\infty} |U(z,t)|^2 dt} \quad (4)$$

*RMS pulse width* in normalized units defined by [4] :

$$s = \left[ \langle t^2 \rangle - \langle t \rangle^2 \right]^{\frac{1}{2}} \quad (5)$$

where :

$$\langle t^n \rangle = \frac{\int_{-\infty}^{\infty} t^n |U(x,t)|^2 dt}{\int_{-\infty}^{\infty} |U(x,t)|^2 dt} \quad (6)$$

RMS pulse width is of interest since it provides a useful metric for assessing performance limitations in fiber-optic communication systems. The RMS pulse width is directly related to the maximum data rate through the commonly used design criterion [1,4] :

$$s_t R_b < \frac{1}{4} \quad (7)$$

where  $s_t$  is *RMS pulse width* at fiber output and  $R_b$  is *bit rate*.

### 2.2 Pre-Emphasis Method

An amplifier is placed periodically along the fiber link to compensate fiber loss and its gain is adjusted so that fiber loss between two amplifiers is exactly compensated.

Fiber loss reduces the effective value of  $N$  as the soliton propagates so in the pre-emphasis method  $N$  is made greater than 1, where :

$$N = \sqrt{\frac{aL}{1 - \exp(-aL)}} \quad (8)$$

Where  $L$  is amplifier spacing and to avoid large fluctuations it is necessary to keep :

$$L \ll \frac{P}{20} L_D \quad (9)$$

### 3 METHODOLOGY

I use SSSF method to solve numerically Eq. (2), algorithm of SSSF method is written using MatLab.

#### 3.1 SSSF Method

Eq. (2) can be written as [2]

$$\frac{\partial U}{\partial x} = \left( \hat{D} + \hat{N} \right) U \quad (10)$$

$\hat{D}$  and  $\hat{N}$  is linear operator and nonlinear operator, where form Eq. (2)

$$\hat{D} = -\text{sgn}(b_2) \frac{i}{2} \frac{\partial^2 u}{\partial t^2} + \text{sgn}(b_3) \frac{1}{6} b_3 \frac{L_D}{L_D} \frac{\partial^3 u}{\partial t^3}; \hat{N} = iN^2 |U|^2.$$

SSSF method nonlinear effects assumed at the middle of step size segment and calculated its effect for the whole segment [3]. The approximation solution of NLSE by SSSF method is :

$$U(x+h, T) = \exp\left(\frac{h}{2} \hat{D}\right) \exp\left(\int_x^{x+h} \hat{N}(x') dx'\right) \exp\left(\frac{h}{2} \hat{D}\right) U(x, T) \quad (11)$$

$$\int_x^{x+h} \hat{N}(x') dx' \approx \frac{h}{2} \left( \hat{N}(x) + \hat{N}(x+h) \right)$$

$h$  is normalized step size parameter. Accuration of SSSF method is third order of  $h$  ( $O(h^3)$ ).  $h$  chosen so that the maximum phase shift ( $\Phi_{max}$ ) due the nonlinear operator is below certain value as

$$g |A_p|^2 h \leq \frac{\Phi_{max}}{L_D} \quad (12)$$

$A_p$  is peak pulse envelope. It has been reported that when  $\Phi_{max}$  is below 0.05 rad, SSSF method gives a good result for simulation of most contemporary optical communication systems [4].

#### 3.2 Soliton Propagation

We simulated soliton propagation with hyperbolic-secant input to the NLSE as

$$U(0, \tau) = \text{sec } h(\tau) \quad (13)$$

First, I compare propagation of fundamental soliton in fiber that loss is zero and in fiber that loss is not zero using pre-emphasis method using Eq. (13) as input to Eq. (2). Then, I use Eq. (13) as input to Eq. (2) to simulate propagation of fundamental soliton with  $\beta_3$  is included and  $\beta_3$  is negligible and compare the result. Finally using rms pulse width, I compare the bit rate transmission distance product of both.

## 4 RESULTS

#### 4.1 Pre-Emphasis Method

Fig.1 is propagation of fundamental soliton in *DSF* as result of *SSSF* using Eq. (13) as input, where  $\alpha = 0$  dB,  $\beta_2 = -1$  ps<sup>2</sup>/km,  $\beta_3 = 0.5$  ps<sup>3</sup>/km,  $T_0 = 1$  ps and  $h = 0.01$ .

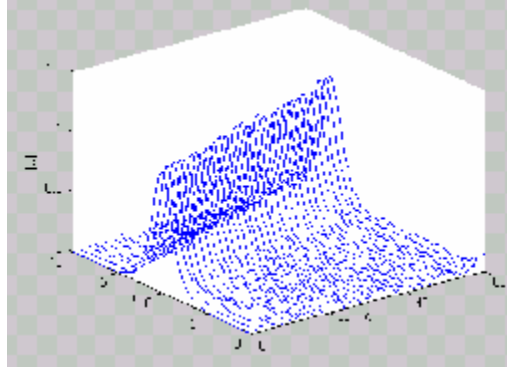


FIGURE 1: Propagation of fundamental soliton,  $\alpha = 0 \text{ dB}$ .

Fig.2 is propagation of fundamental soliton in *DSF* as result of *SSFM* and pre-emphasis method with  $L = 0,1L_D$ , using Eq. (13) as input, where  $\alpha = 0.2 \text{ dB}$ ,  $\beta_2 = -1 \text{ ps}^2/\text{km}$ ,  $\beta_3 = 0.5 \text{ ps}^3/\text{km}$ ,  $T_0 = 1 \text{ ps}$  and  $h = 0.01$ .

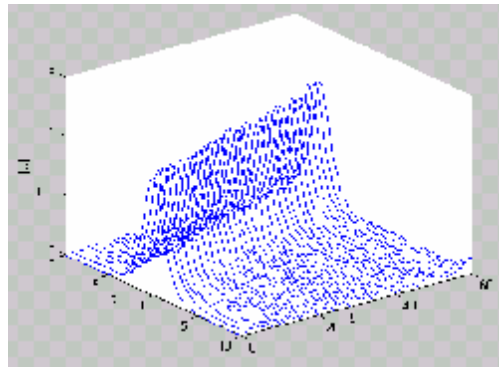
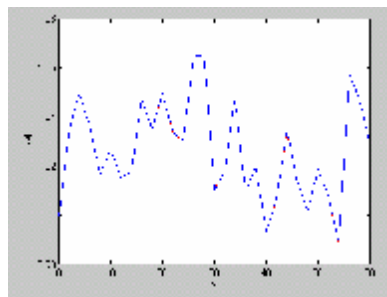
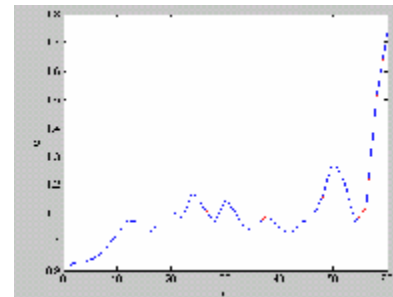


FIGURE 2 : Propagation of fundamental soliton,  $\alpha = 0.2 \text{ dB}$ .

Fig. 3(a) and 3(b) is amplitude comparison and RMS pulse width comparison for  $\alpha = 0 \text{ dB}$  and  $\alpha = 0.2 \text{ dB}$  as function of propagation distance. Fig. 3(c) is comparison of them, using their absolute difference values.



(a)



(b)

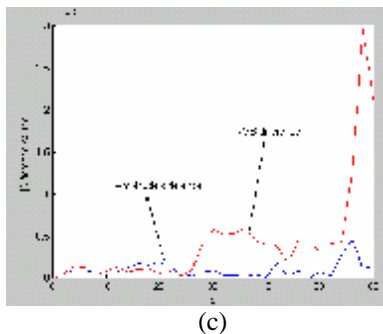


FIGURE 3 : (a) Amplitude comparison. (b) RMS pulse width comparison. (c) Amplitude and RMS pulse width difference.

#### 4.2 Higher Order Dispersion Effects

Fig.4 is propagation of fundamental soliton in *DSF* as result of *SSSFM* and pre-emphasis method with  $L = 0,1L_D$  using Eq. (13) as input with  $\beta_3$  is negligible ( $\beta_3 = 0$ ), where  $\alpha = 0.2$  dB,  $\beta_2 = -1$  ps<sup>2</sup>/km,  $T_0 = 1$  ps and  $h = 0.01$ . While, at Fig.2 we can see propagation of fundamental soliton with  $\beta_3$  is included.

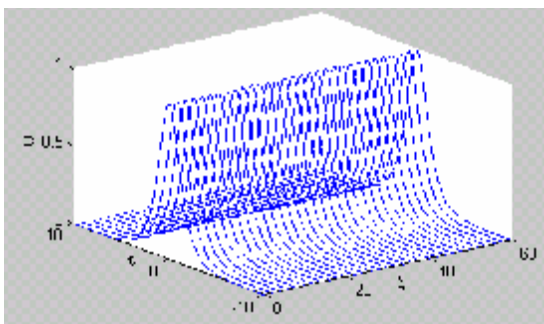
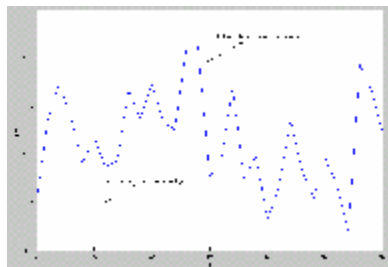
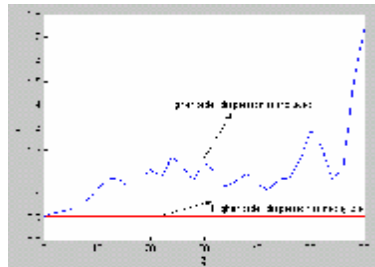


FIGURE 4 : Propagation of fundamental soliton,  $\beta_3=0$ .

Fig.5 (a) and 5 (b) is comparison of amplitude and RMS-pulse width of fundamental soliton propagation in *DSF* as result of with  $\beta_3$  is negligible ( $\beta_3 = 0$  ps<sup>3</sup>/km) and  $\beta_3 = 0.5$  ps<sup>3</sup>/km. Fig. 5(c) is comparison of them, using their absolute difference values [4].



(a)



(b)

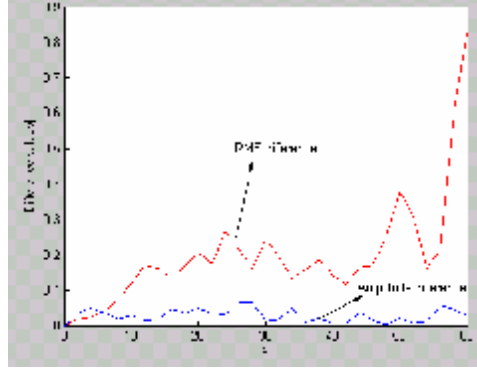


FIGURE 5 : (a) Amplitude comparison. (b) RMS pulse width comparison. (c) Amplitude and RMS pulse width difference.

## 5 DISSCUSION

We can compute that average value of amplitude difference between results of SSSFm with  $\alpha = 0$  and  $\alpha = 0.2$  dB that compensated by pre-emphasis method along propagation ( $\zeta = 60$ ) is  $9.6096 \cdot 10^{-5}$  and  $4.3362 \cdot 10^{-4}$  for rms pulse width difference. So we can conclude that pre-emphasis method gives a accurate result to compensate fiber loss.

From Fig.(2) and Fig.(4) we can see that effects of higher dispersion change the position of pulse peak so the pulse is not symmetry anymore around  $\tau = 0$  and make ripples near leading and trailing edge of the pulses.

Table1 is rms pulse width as function of normalized distance from Fig 3.(a) and 3.(b), for  $\zeta = 0, 10, 20, 30, 40, 50$  and  $60$ .

TABLE 1 : RMS pulse width as function of  $\zeta$ .

$\zeta$ (Normalized distance)	$\sigma$ (for $\beta_3 = 0$ ps <sup>3</sup> /km)	$\sigma$ (for $\beta_3 = 0.5$ ps <sup>3</sup> /km)
0	0.90690	0.9069
10	0.90700	1.028
20	0.90712	1.1115
30	0.90708	1.1497
40	0.90695	1.0484
50	0.90694	1.2873
60	0.90703	1.7439

From table1 we can see that for fiber optics with  $\beta_3 = 0$  so the rms pulse width is almost constant, using Eq. (7) we get constant bit rate of transmission is about,  $R_b < 0.2756$  Gb/s . Finally we use the value of rms pulse width in table1 and compute bit rate transmission distance product, the result is in table2.

TABLE 2 : Bit rate transmission distance product.

$L_T$ (Km) = $\zeta \cdot L_D$ (Distance)	$R_b L_T$ (Gb/s.Km) (for $\beta_3 = 0$ ps <sup>3</sup> /km)	$R_b L_T$ (Gb/s.Km) (for $\beta_3 = 0.5$ ps <sup>3</sup> /km)
10	< 2.756	< 2.431907
20	< 5.512	< 4.498426
30	< 8.268	< 6.523441
40	< 11.024	< 9.538344
50	< 13.78	< 9.710246
60	< 16.536	< 8.601411

From table2 we can conclude that effects of higher order dispersion is also reduce the bit rate transmission distance product.

## 6 CONCLUSION

In his paper, the simulation results show that pre-emphasis method gives a accurate result to compensate fiber loss.

When higher order dispersion is negligible, the pulse shape (peak and rms pulse width) is almost not change along the propagation so give a almost constant bit rate for the system. When higher order dispersion is included so it's effects will increase the rms pulse width of pulse along the propagation which reduce the bit rate transmission distance product.

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